

Application of vector analysis to alternating current ckt (1)

Introduction:- Representation of instantaneous value of alternating current as a scalar product.

Let \vec{I} be a vector of constant magnitude \hat{I} , rotating in positive (counter-clockwise) sense with constant angular velocity ' ω ', its slope, relative to unit vector \vec{v} , at time $t=0$ being ' ϕ '.

' i ' be instantaneous current $\hat{I} \cos(\omega t + \phi)$.

Since at time ' t ' the vector \vec{I} makes an angle $(\omega t + \phi)$ with \vec{v} .

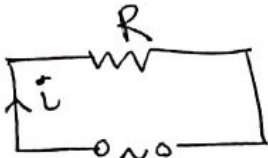
$$\therefore i = \vec{I} \cdot \vec{v} = \hat{I} \cos(\omega t + \phi) \vec{v} \cdot \vec{v}$$

\therefore An alternating current of amplitude \hat{I} and angular freq ω can be represented as a scalar product with \vec{v} of a vector \vec{I} , ~~makes an angle $(\omega t + \phi)$ with \vec{v}~~ , of constant magnitude \hat{I} , rotating with angular velocity ' ω ', initial position of \vec{I} making an angle ϕ with \vec{v} .

Current and potential relationship in vector form

a) Pure resistance (R)

Suppose alternating potential difference $\vec{V} \cos(\omega t)$ applied to non-inductive resistance R



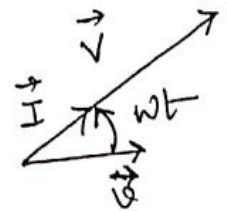
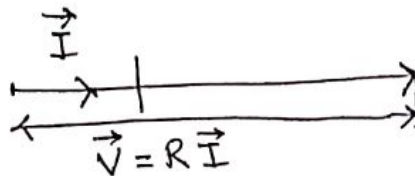
Then ' i ' \rightarrow instantaneous value of current

$$Ri = V$$

$$\text{or } R \vec{I} \cdot \vec{v} = \vec{V} \cdot \vec{v} \quad | \text{ taking projection along } \vec{v}$$

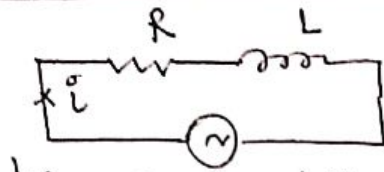
$$\therefore \vec{V} = R \vec{I}$$

Voltage and current are in phase



Inductive Resistance (R and L in series) :-

(2)



Let, $v = \hat{v} \cos(\omega t)$
 $= \hat{v} \exp(j\omega t)$
 $\vec{v} \cdot \vec{v}$
 $i = \vec{I} \cdot \vec{v}$

Now, $v = iR + L \left(\frac{di}{dt} \right)$
 $\vec{v} \cdot \vec{v} = \hat{v} \exp(j\omega t) \vec{v} \cdot \vec{v}$
 $= R \vec{I} \cdot \vec{v} + j\omega L \vec{I} \cdot \vec{v}$
 $\therefore \vec{I} = \frac{\hat{v} \exp(j\omega t) \vec{v}}{(R + j\omega L)}$
 $= \frac{\vec{v}}{Z}$

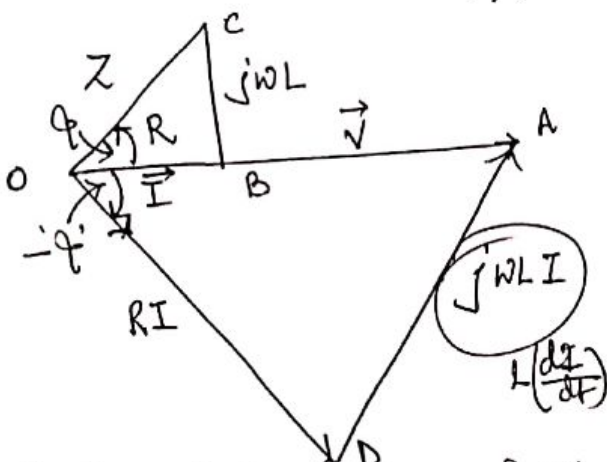
When, $Z = R + j\omega L$ is known as vector impedance or impedance operator.

Let, $\frac{1}{(R + j\omega L)}$
 $= \frac{R - j\omega L}{(R^2 + \omega^2 L^2)} = \alpha - j\beta = \rho \exp(-j\phi)$

$\alpha = \frac{R}{(R^2 + \omega^2 L^2)}$, $\beta = \frac{\omega L}{(R^2 + \omega^2 L^2)}$. $\tan \phi = \left(\frac{\omega L}{R} \right)$

$\therefore \vec{I} = \frac{\hat{v} \exp(j\omega t) \vec{v}}{(R + j\omega L)}$
 $= \frac{\hat{v}}{\sqrt{R^2 + \omega^2 L^2}} \exp(j\omega t) \cdot \exp[-j\phi] \vec{v}$
 $= \frac{\hat{v}}{\sqrt{R^2 + \omega^2 L^2}} \exp[j(\omega t - \phi)] \vec{v}$

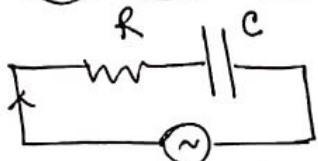
$\therefore i = \vec{I} \cdot \vec{v} = \hat{I} \cos(\omega t - \phi)$



Here \vec{I} is a vector of magnitude $\hat{I} = \frac{\hat{v}}{\sqrt{R^2 + \omega^2 L^2}}$ rotating with angular velocity ω , and whose angular position at $t=0$ is, $-\tan^{-1}\left(\frac{\omega L}{R}\right)$ i.e. phase of current is always behind impressed voltage by amount ' ϕ '.

' ωL ' is inductive reactance of coil. \vec{OB} horizontally equal to ' R ', ' BC ' normal to AB
 \rightarrow length ωL is ' $j\omega L$ '. $\Rightarrow \vec{OD} + \vec{DA} = \vec{OA} = \vec{v}$ $\vec{OD} = \vec{IR}$, $\vec{DA} = j\omega L \vec{I}$
 \vec{OB} and \vec{OD} are voltage vectors gives potential difference betⁿ R and L

Imperfect Condenser i.e R and C in series



Imperfect Condenser may be regarded as Capacitance 'C' in series with resistance 'R'

$V = \hat{V} \cos(\omega t)$

$V = Ri + \frac{q}{C}$
 $= Ri + \frac{1}{C} \int i dt$

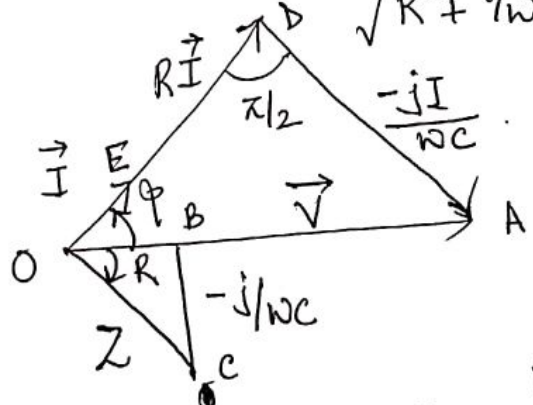
$\vec{V} \cdot \vec{V} = R \vec{I} \cdot \vec{V} + \frac{1}{j\omega C} \vec{I} \cdot \vec{V}$
 $\therefore \vec{V} = R \vec{I} + \frac{1}{j\omega C} \vec{I}$
 $= (R - \frac{j}{\omega C}) \vec{I}$

as $i = \hat{I} \sin(\omega t) \vec{V}$
 $\int i dt = \frac{\hat{I}}{j\omega} \sin(\omega t) \vec{V}$

$\therefore \vec{I} = \frac{\vec{V}}{(R - j/\omega C)} = \frac{(R + j/\omega C)}{(R^2 + 1/\omega^2 C^2)} \vec{V}$
 $= \frac{\hat{V}}{\sqrt{R^2 + 1/\omega^2 C^2}} [\sin(\omega t + \phi)] \vec{V}$

when $\vec{V} = \hat{V} e^{j\omega t}$
 and $\tan \phi = \frac{1}{\omega RC}$

$\therefore i = \vec{I} \cdot \vec{V} = \frac{\hat{V}}{\sqrt{R^2 + 1/\omega^2 C^2}} \cos(\omega t + \phi)$



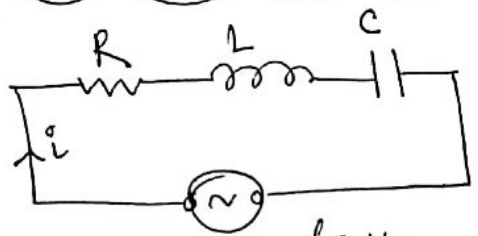
Horizontal to represent \vec{V} .
 $OB = R$, $BC = -j/\omega C$ (vertically downward)
 $OC = Z = (R - j/\omega C)$

Also, $\vec{I} = \frac{\vec{V}}{Z}$, $OE = \frac{\hat{V}}{\sqrt{R^2 + 1/\omega^2 C^2}}$ and it is advance of ϕ of OA

$OD = R \vec{I}$, $DA = -j/\omega C \vec{I}$

Resistance, inductance and Capacitance in Series: -

(4)



$$v = iR + \frac{\int i dt}{c} + L \left(\frac{di}{dt} \right)$$

$$\therefore v = \vec{v} \cdot \vec{v} = \hat{v} (\cos(j\omega t)) \vec{v} \cdot \vec{v}$$

$i = \vec{I} \cdot \vec{v}$, \vec{I} has to be determined.

$$\hat{v} [\cos(j\omega t)] \vec{v} \cdot \vec{v} = \left(R + j\omega L + \frac{1}{j\omega c} \right) \vec{I} \cdot \vec{v}$$

$$\therefore \vec{I} = \frac{\hat{v}}{R + j(\omega L - \frac{1}{\omega c})} (e^{j\omega t}) \vec{v}$$

$$= \frac{\left\{ R - j(\omega L - \frac{1}{\omega c}) \right\} \hat{v} \cos(j\omega t) \vec{v}}{\left[R^2 + (\omega L - \frac{1}{\omega c})^2 \right]}$$

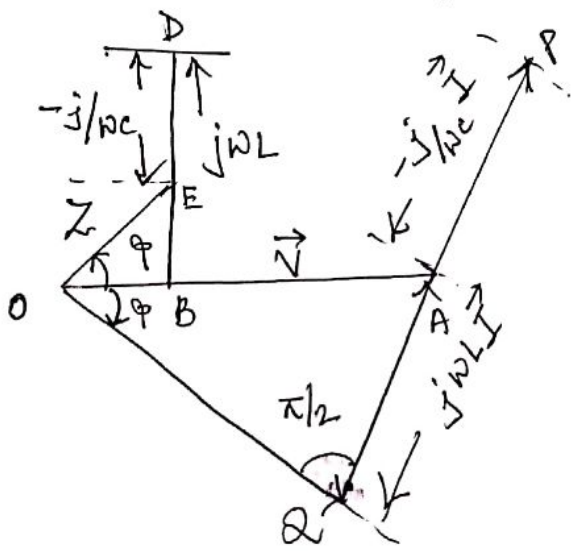
$$\therefore \vec{I} = \frac{\hat{v}}{\sqrt{R^2 + (\omega L - \frac{1}{\omega c})^2}} [\cos(j\omega t - \phi)] \vec{v}$$

$$\phi = \tan^{-1} \left(\frac{\omega L - \frac{1}{\omega c}}{R} \right)$$

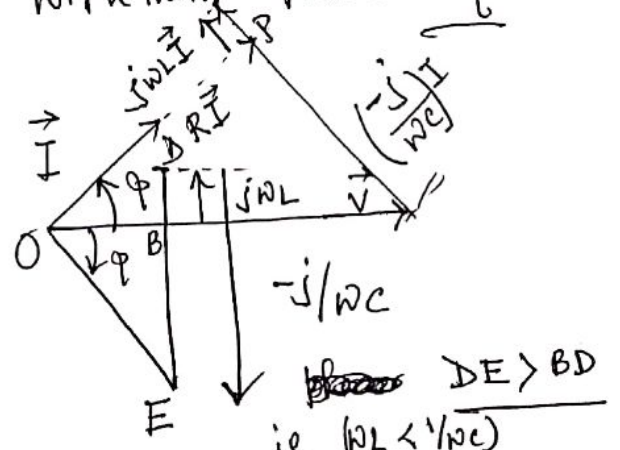
Thus \vec{I} is a vector of magnitude

$$\hat{I} = \frac{\hat{v}}{\sqrt{R^2 + (\omega L - \frac{1}{\omega c})^2}}$$

rotating with angular velocity ' ω ' with initial phase ' $-\phi$ '

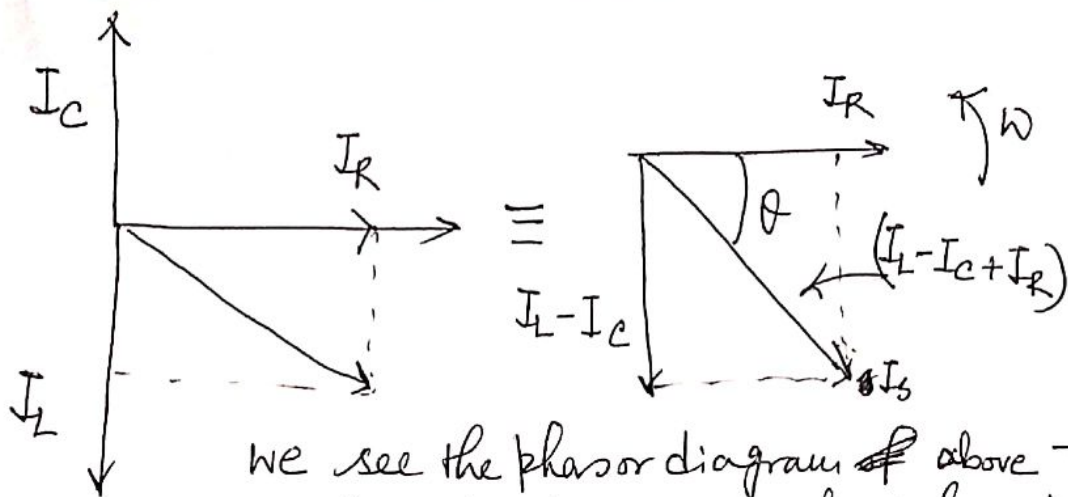


If $(\omega L) > (1/\omega c)$
phase of current is always behind impressed voltage



~~DE > BD~~
ie $(\omega L < 1/\omega c)$
ie ϕ is -ve, current leads voltage.

Phasor Diagram of Parallel LCR CRT



We see the phasor diagram above that the current vectors produce a rectangular triangle, comprising of hypotenuse I_S , horizontal axis I_R and vertical axis $(I_L - I_C)$. It forms a Current triangle

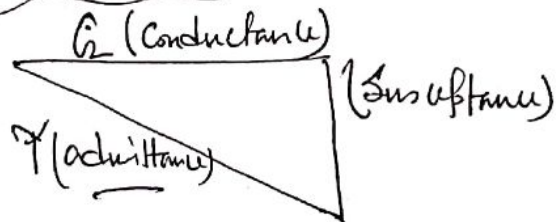
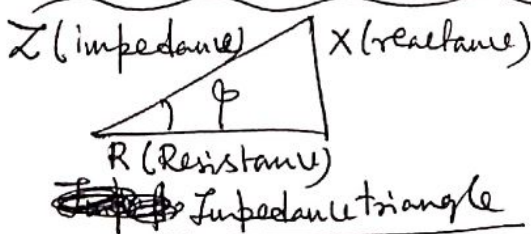
We can use pythagoras's theorem on this current triangle to mathematically obtain the magnitude of branch currents along x-axis and y-axis. then determine total current I_S of these components as shown.

Current Triangle for Parallel RLC CRT :-

$$I_S^2 = I_R^2 + (I_L - I_C)^2$$

$$\therefore I_S = \sqrt{\left(\frac{V}{R}\right)^2 + \left(\frac{V}{X_L} - \frac{V}{X_C}\right)^2} = \frac{V}{Z}$$

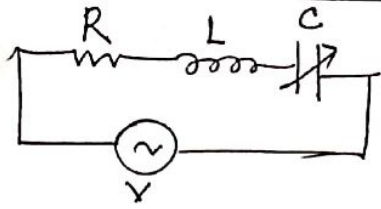
Admittance triangle for parallel RLC CRT :-



$$Z = \frac{1}{Y} = \frac{1}{\sqrt{\left(\frac{1}{R}\right)^2 + \left(\frac{1}{\omega L} - \omega C\right)^2}}$$

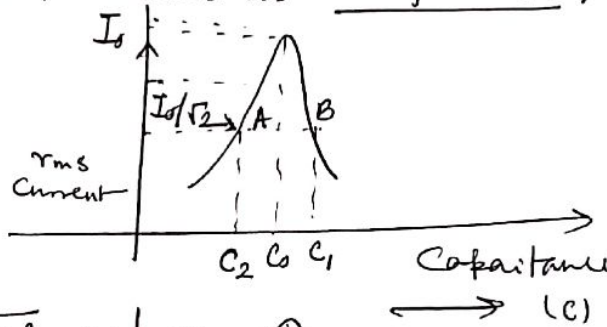
Series-Resonance (Current Resonance) :-

(5)



'v' be rms potential difference applied to simple series resonance ckt. ckt is known as acceptor ckt, (with variable C)

$$I^2 = \frac{V^2}{R^2 + (\omega L - 1/\omega C)^2}$$

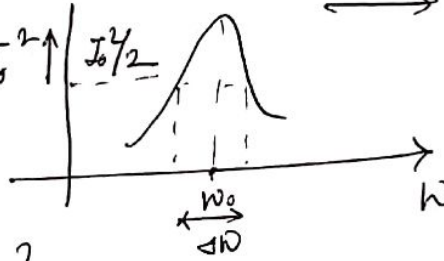


$$\therefore I^2 = \frac{V^2}{R^2} \times \frac{1}{1 + \frac{1}{R^2} \left[\frac{1}{\omega_0 C_0} - \frac{1}{\omega_0 C} \right]^2}$$

at resonance

$$\omega_0 L = \frac{1}{\omega_0 C_0} \Rightarrow \omega_0 = \frac{1}{\sqrt{LC_0}}$$

$$I_0 = V/R \rightarrow \text{Current at resonance}$$



$$\therefore \left(\frac{I_0}{I} \right)^2 - 1 = \frac{1}{\omega_0^2 R^2} \left[\frac{1}{C_0} - \frac{1}{C} \right]^2$$

choosing, $(I_0/I)^2 = 2$

$$\therefore \omega_0^2 R^2 = \left(\frac{1}{C_0} - \frac{1}{C} \right)^2 = \left(\frac{C - C_0}{C C_0} \right)^2$$

$$\therefore \omega_0 R = \frac{(C - C_0)}{C C_0} \text{ or } \frac{(C_0 - C_2)}{C_0 C_2}$$

when, C_1 and C_2 are values of C

and, $1/C_1 + 1/C_2 = 2/C_0$

In practice, C_1 and C_2 don't differ very much from C_0

$$\therefore C_0 C_1 = C_0 C_2 = C_0^2$$

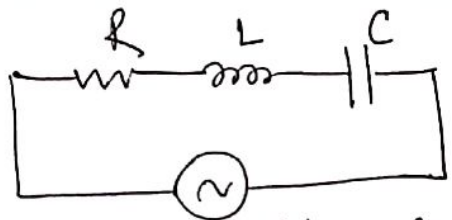
$$\therefore \frac{C_1 - C_2}{C_0^2} = 2 \omega_0 R$$

$$\rightarrow \frac{(C_1 - C_2)}{2 C_0} = \omega_0 C_0 R = \frac{R}{\omega_0 L}$$

$\frac{2 C_0}{C_1 - C_2}$ is measure of peaking of response curve i.e. selectivity of ckt. denoted by 'Q'

$$\therefore \boxed{Q = \text{Quality factor / Magnification} = \frac{\omega_0 L}{R}} \checkmark$$

Series Resonance (Current resonance)



$$I = \frac{V}{\sqrt{R^2 + (\omega L - \frac{1}{\omega C})^2}} = \frac{V/L}{\sqrt{(\omega - \frac{1}{\omega LC})^2 + R^2/L^2}}$$

Now, $\omega_0 = \frac{1}{\sqrt{LC}}$

$$\text{or, } I^2 = \frac{V^2/L^2 \omega^2}{\sqrt{(\omega_0^2 - \omega^2)^2 + \alpha^2 \omega^2}}$$

$$\leftarrow \therefore I = \frac{V/L}{\sqrt{\frac{(\omega^2 - \omega_0^2)^2}{\omega^2} + R^2/L^2}}$$

$$\therefore R^2/L^2 = \alpha^2$$

If I_0 be current at resonance,

$$\therefore I_0 = \sqrt{\frac{V^2/L^2 \cdot \omega_0^2}{R^2/L^2 \omega_0^2}} = \frac{V}{R}$$

In order to have, $I_0^2/I^2 = 2$, let the frequency of the supply be such that, $\omega = \omega_0 \pm \frac{1}{2} \Delta\omega$

$$(\omega - \omega_0)^2 = \Delta\omega^2$$

$$\begin{aligned} I^2 &= \frac{V^2/L^2 \cdot \omega_0^2}{\alpha^2 (\omega_0 \pm \frac{1}{2} \Delta\omega)^2 + (\omega_0 + \omega)^2 (\omega_0 - \omega)^2} \\ &= \frac{V^2/L^2 [\omega_0 \pm \frac{1}{2} \Delta\omega]^2}{\alpha^2 (\omega_0 \pm \frac{\Delta\omega}{2})^2 + (2\omega_0 \pm \frac{1}{2} \Delta\omega)^2 \cdot \frac{1}{4} (\Delta\omega)^2} \end{aligned}$$

$$\begin{aligned} I^2 &= \frac{(V^2/L^2) (\omega_0 \pm \frac{1}{2} \Delta\omega)^2}{[\omega_0^2 - (\omega_0 \pm \frac{1}{2} \Delta\omega)^2]^2 + \alpha^2 [\omega_0 \pm \frac{1}{2} \Delta\omega]^2} \\ &= \frac{V^2/L^2 (\omega_0 \pm \frac{1}{2} \Delta\omega)^2}{\left[-\frac{(\Delta\omega)^2}{4} \mp \omega_0 \Delta\omega\right]^2 + \alpha^2 \omega_0^2 \pm \alpha^2 \omega_0 \Delta\omega + \frac{\alpha^2 \Delta\omega^2}{4}} \end{aligned}$$